

Why Photon Colliders are necessary in a future collider program

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Ilya F. Ginzburg

Sobolev Institute of Mathematics SB RAS, 630090, Novosibirsk, Russia

Subjects of discussion

1. Different scenarios
2. Higgs window to a New Physics
3. Anomalies in the interactions of gauge and Higgs bosons
4. Some problems with t -quarks and \mathcal{MSSM} .
5. Some problems in QCD and hadron physics.
6. By-product: Production of axions, etc. from region of conversion $e \rightarrow \gamma$.

1 Different scenarios

Photon Colliders in the widespread scenario. When discussing the program for future high energy colliders, the basic point is usually that *Nature is so favorable to us to dispose the essential fraction of new particles and the thresholds of new interactions* (e.g., effects of new higher dimensions, compositeness, etc.) *within the LHC operation domain*.

In this case main discoveries will be made at the Tevatron and the LHC. The e^+e^- Linear Colliders (LC) in their first stages will be machines for measuring precise values of coupling constants and exploring in detail supersymmetry. The high luminosity expected for e^+e^- Linear Colliders provides opportunity to obtain parameters of models realized with very high accuracy, see, e.g., [1].

With the Photon Collider mode of LC [2] having roughly the same luminosity as that for e^+e^- mode [3] these results will be improved. Indeed,

- The cross sections of production of pairs of charged particles in the $\gamma\gamma$ collision are $5 \div 8$ times higher than that in the e^+e^- collision in its maxima; this ratio increases with growth of energy.
- The set of final states at a Photon Collider is much richer than that at in an e^+e^- mode.
- One can vary polarizations of photon beams relatively easily (*circular* \rightarrow *transverse*).

Unfortunately, this Photon Collider opportunity was not explored by the community in necessary detail. It can be assured that the forthcoming studies of many physicists

with detailed simulation will show an exceptional potential for the Photon Collider mode in both new problems and all problems considered till now.

Possible \mathcal{SM} – like scenario. It can also happen that the opposite scenario will be realized:

No new particles and interactions will be discovered at the Tevatron, LHC and e^+e^- LC, except the Higgs boson. If additionally its coupling constant to quarks and gauge bosons will be close to their values in the \mathcal{SM} , we find a \mathcal{SM} – like picture of the World. It can be realized both if our World is really described by simple \mathcal{SM} up to very small distances and if some other model is realized and New Physics is *round the corner*. In this case the main goal of studies at new colliders will be the hunting for indirect signals of New Physics – deviations of observed quantities from \mathcal{SM} predictions. *Photon Colliders are the best machines for the hunting for signals of New Physics if the \mathcal{SM} – like scenario is realized.*

2 Higgs window to a New Physics

The study of Higgs-boson couplings with photons ($h\gamma\gamma$ and $hZ\gamma$) looks like a very promising tool for resolving of models of New Physics.

- These couplings are absent in the \mathcal{SM} at tree level, appearing only at the loop level. Therefore, the background for signals of New Physics will be relatively lower here than in other processes which are allowed at tree level of the SM.
- All fundamental charged particles contribute to these effective couplings. The whole structure of the theory influences the corresponding Higgs-boson decays.
- The expected accuracy in the two-photon width is *about 2%* at $M_h \leq 150$ GeV and even at the luminosity integral $\sim 30 \text{ fb}^{-1}$ in the high energy peak of the luminosity spectrum [4].

In the $2\mathcal{HDM}$ and \mathcal{MSSM} , the observed Higgs boson will be either the lightest Higgs boson h or heavier one H . Assuming the coupling constants of observed Higgs boson to quarks and gauge bosons are close to their values in the \mathcal{SM} , other Higgs bosons can easily avoid observation at LHC and e^+e^- LC due to small couplings to standard matter. These neutral Higgses can be seen in the process $\gamma\gamma \rightarrow h$ [5] and charged Higgses in the process $\gamma\gamma \rightarrow H^+H^-$.

Even if these additional Higgs bosons are so heavy that they cannot be produced at a first stage Photon Collider, the models can be distinguished well via precise measurement of two-photon width of observed \mathcal{SM} – like Higgs boson [6]. In this paper we assumed that the discussed \mathcal{SM} – like scenario is realized with an accuracy estimated for \mathcal{SM} Higgs boson at an e^+e^- Linear Collider [1]. The two photon width is calculated via the Higgs couplings with the matter measured at an e^+e^- LC. In the $2\mathcal{HDM}$, deviation from \mathcal{SM} is due mainly to contribution of charged Higgs. In the general case the two-photon width is about 10% less than that in the \mathcal{SM} (see Figs. 1). It is several times larger than the expected experimental inaccuracy. Therefore, measurement of the two-photon width at a Photon collider can resolve these models reliably.

In the \mathcal{MSSM} the lightest Higgs boson is decoupled with superpartners and other members of Higgs multiplet if they are very heavy. More detailed calculations will give us

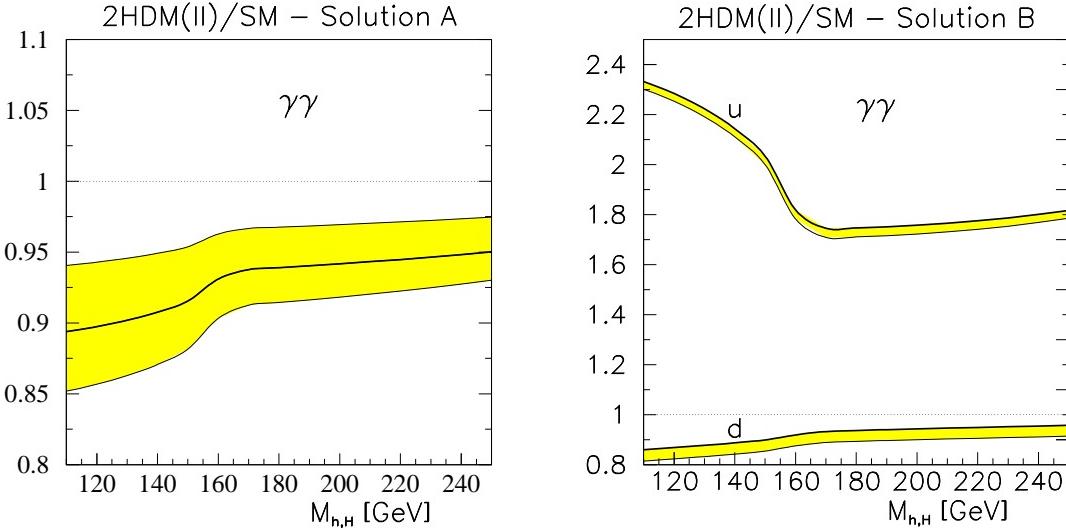


Figure 1: *The ratio of two-photon width of Higgs boson in $2\mathcal{HDM}$ to its \mathcal{SM} value at $\lambda_5 = 0$ and $M_h \geq 800$ GeV. Shaded zones correspond to experimental uncertainties expected in the experiments at e^+e^- LC. Deviation from \mathcal{SM} depends on λ_5 as $\propto (1 - \lambda_5 v^2/M_{H^\pm}^2)$ with $v = 246$ GeV is v.e.v. of Higgs field. Solution A is really close to \mathcal{SM} . For the solution B some of the couplings of the Higgs field with matter are close to the \mathcal{SM} values, the others have the opposite sign.*

the upper bounds for masses of superparticles which can influence the photon widths so that the difference will be seen in the experiment. Preliminary estimates show that these values are higher than the discovery limits at LHC.

In many variants of $2\mathcal{HDM}$ or $MSSM$ the masses of heavy scalar Higgs H and Higgs pseudoscalar A are close each other. In some other variants they are mixed (\mathcal{CP} violated scenarios). The observations at LHC and e^+e^- Linear Collider cannot often resolve these opportunities due to low resolution for these bosons. The polarization asymmetries in Higgs boson production at a Photon Collider can resolve these variants, i.e. establish whether \mathcal{CP} parity at Higgs level is violated or not.

3 Anomalies in the interactions of gauge and Higgs bosons

Before discovery new heavy particles inherent New Physics, it reveals itself at lower energies as some anomalies in the interactions of known particles. Our goal is to find these anomalies and discriminate them as best as possible. *The correlation between coefficients of different anomalies will be the key for understanding what is the nature of New Physics.*

Interactions of gauge bosons. The practically unique process under study in the e^+e^- mode is $e^+e^- \rightarrow W^+W^-$. At suitable electron polarization the neutrino exchange contribution (having small interest) disappears, and the residual cross section (obliged by photon and Z boson exchange) has a maximum about 2 pb within the LEP operation

interval and decreases with energy after that. The cross sections of other processes with W production ($e^+e^- \rightarrow e^+e^- WW$, $e^+e^- \rightarrow e\nu W$, ...) are small at $\sqrt{s} < 1$ TeV.

At Photon collider the main processes are $\gamma\gamma \rightarrow W^+W^-$, $e\gamma \rightarrow \nu W$. Their cross sections are about 80 pb and are energy independent at $\sqrt{s} > 200$ GeV. That is at least 40 times higher than for e^+e^- collisions, it gives about $(1 \div 5) \cdot 10^7$ W 's per year, comparable with Z production at LEP. Due to high value of these basic cross sections, many processes of 3-rd and 4-th order have large enough cross sections:

$\gamma e \rightarrow eWW$, $\gamma\gamma \rightarrow ZWW$, $\gamma e \rightarrow \nu WZ$, $\gamma\gamma \rightarrow WWW$, $\gamma\gamma \rightarrow WWZZ$, etc. Large variety of these processes permit us to discover and separate well anomalies in specific processes and (or) distributions. This subject needs studies of all processes enumerated with details of behavior in different regions of phase space and polarization dependence.

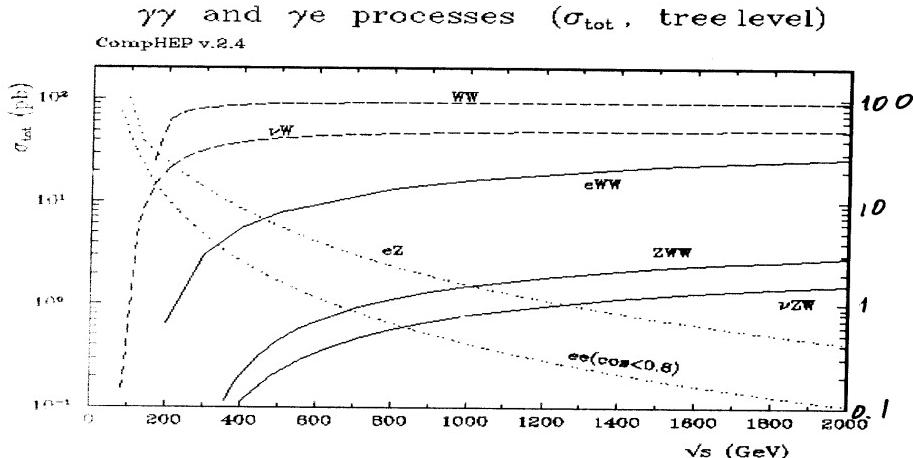


Figure 2: Some processes of gauge boson production at γe and $\gamma\gamma$ colliders

- For example, the following line of procedures is almost evident:
 - To extract γWW anomalies from $\gamma e \rightarrow \nu W$.
 - To extract ZWW anomalies from $e^+e^- \rightarrow W^+W^-$.
 - To extract $\gamma\gamma WW$ anomaly from $\gamma\gamma \rightarrow W^+W^-$.
- $e\gamma \rightarrow \nu W$. The cross section of process is $\propto (1 - 2\lambda_e)$, it is switched on or off with variation of electron helicity λ_e . It gives precise test of absence of right handed currents in the interaction of W with the matter.
- $\gamma e \rightarrow eWW$. The cross section of process at $\sqrt{s} = 500$ GeV is about 10 pb, corresponding counting rate is few millions events per year.
- ★ In events with transverse momentum of scattered electron $p_\perp \geq 30$ GeV one can study here anomaly γZWW .
- ★ The charge asymmetry of produced W 's looks like most sensitive key for the study of strong interaction in Higgs sector.
- $\gamma\gamma \rightarrow W^+W^-$ and $e\gamma \rightarrow \nu W$. The two-loop radiative corrections to these processes are measurable and sensitive to the problems:
 - ★ construction of S -matrix of theory with unstable particles;
 - ★ gluon corrections like Pomeron exchange between quark components of W 's.

Interactions of Higgs boson with light ($\gamma\gamma \rightarrow h$, $\gamma e \rightarrow eh$). All observable anomalous \mathcal{CP} – even and \mathcal{CP} – odd interactions of Higgs boson with light are summarized in the form of an effective Lagrangian:

$$\Delta\mathcal{L} = 2Hv \left(\theta_\gamma \frac{F_{\mu\nu}F^{\mu\nu}}{2\Lambda_\gamma^2} + \theta_Z \frac{Z_{\mu\nu}F^{\mu\nu}}{\Lambda_Z^2} + i\theta_{P\gamma} \frac{F_{\mu\nu}\tilde{F}^{\mu\nu}}{2\Lambda_{P\gamma}^2} + i\theta_{PZ} \frac{Z_{\mu\nu}\tilde{F}^{\mu\nu}}{\Lambda_{PZ}^2} \right), \quad (\theta_i = e^{i\xi_i}). \quad (1)$$

Here $F^{\mu\nu}$ and $Z^{\mu\nu}$ are the standard field strengths for the electromagnetic and Z field, $\tilde{F}^{\mu\nu} = \epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta}/2$ and ξ_i are the phases of couplings, generally different from 0 or π even in the C–even case (due to possible anomalous contributions of some light particles) [7].

While \mathcal{CP} even anomalies can be seen via the values of measured cross sections, the \mathcal{CP} anomalies will be seen via the polarization asymmetries ("longitudinal" – variation of cross sections with change of sign of collided photon helicities, or "transverse" – variation of cross section in dependence on the angle between directions of linear polarization of collided photons). The effects are large enough to see them at reasonable values of anomaly scales Λ_i and phases ξ_i (see Fig. 2).

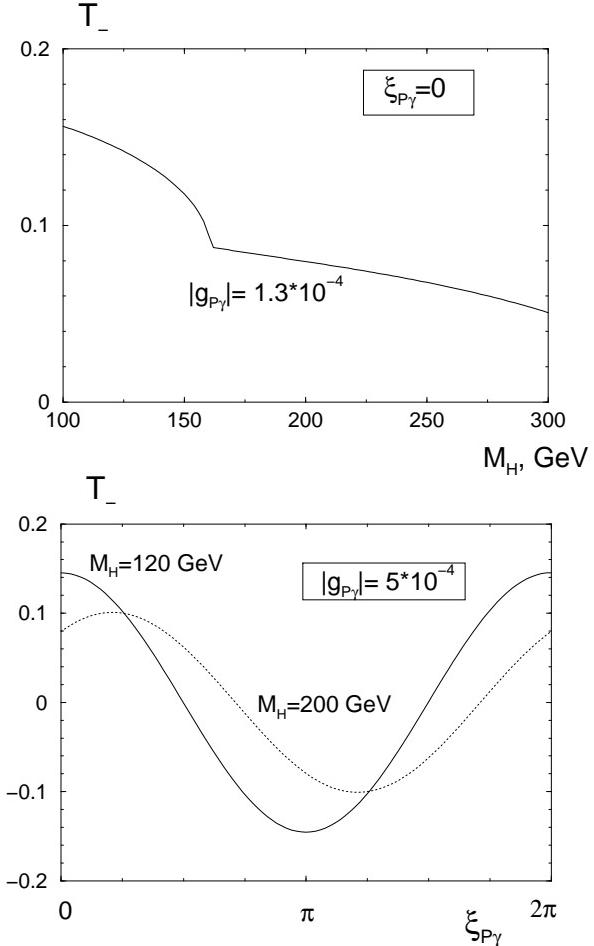


Figure 3: "Longitudinal" asymmetry for $\gamma\gamma \rightarrow H$ process

The specific example of anomaly presents \mathcal{CP} violating mixing in Higgs sector of

$2\mathcal{HDM}$ (scalar – pseudoscalar) with mixing angle α_2 . We present results in Fig. 3 fixing other model parameters to be close to the \mathcal{SM} case.

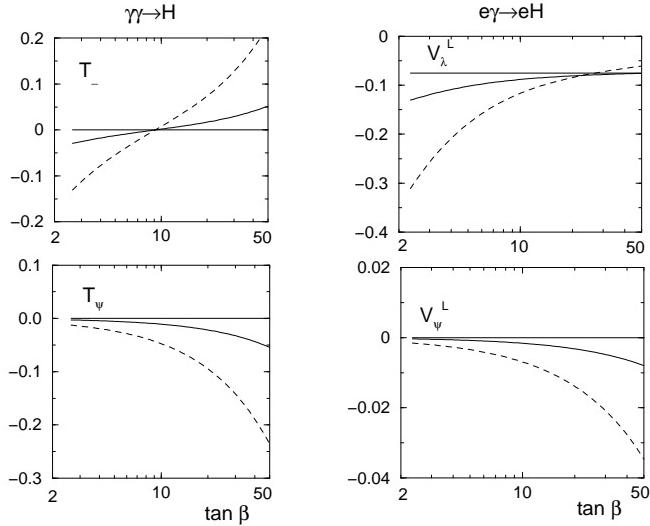


Figure 4: *Polarization asymmetries in $\gamma\gamma \rightarrow h$ and $e\gamma \rightarrow eh$ processes due to scalar-pseudoscalar mixing in $2\mathcal{HDM}$ (II), $M_h = 110$ GeV; for $e\gamma \rightarrow eH$ process $E_{tot} = 1.5$ TeV. The thick solid lines show the \mathcal{SM} values; thin solid and dashed lines refer to $\sin \alpha_2 = 0.1$ and 0.5 respectively (α_2 is light scalar–pseudoscalar mixing angle).*

- It can happen that the masses of heavy scalar Higgs H and pseudoscalar A are so close to each other that they cannot be resolved in the mass spectrum. The study of decay products cannot distinguish this overlapping from the case of their mixing with \mathcal{CP} violation. In the case when these Higgses overlap in the true mass spectrum, the correlations in the decay products will be identical if \mathcal{CP} violated mixing exists or not. The study of polarization asymmetries in the Higgs boson production at Photon Colliders distinguish these two opportunities well.

4 t – quarks, \mathcal{MSM}

Here I present only few examples in which Photon Colliders look absolutely necessary.

- **t -quarks.** In addition to the usually discussed problems related to t -quarks, the specific one is the study of axial anomaly in the process $ye \rightarrow ett\bar{t}$. It was found that at small transverse momenta of scattered electrons p_\perp the cross section of subprocess $Z_L\gamma \rightarrow tt$ with longitudinally polarized Z does not disappear as happens with photons but diverges as M_t^2/p_\perp^2 [8]

• Some problems with \mathcal{MSM} .

- ★ Due to high values of the basic cross sections for the production of pairs of charged particles, Photon Colliders promise to be an excellent place for observation of (even small) possible flavor nonconservation in the neutral currents with superparticles.
- ★ If the stop squark is not very heavy, the atom-like stoponium with mass 200-400 GeV

should also exists. It will be very narrow and hard to observe at hadron collider and e^+e^- Linear Colliders. However, it can be clearly seen at $\gamma\gamma$ collider with cross section averaged over photon spectrum $\langle \sigma \rangle \approx 10 - 50 \text{ fb}$ and clear enough signature [9].

* Gaugino discovery. The expected discovery limit for gaugino at LHC is about 170 GeV in the variants of MSSM which looks most probable now. It can be enhanced till 450 GeV in the "less probable" variants of MSSM.

5 QCD and Hadron Physics

All problems studied at HERA and LEP will be studied at Photon Colliders but in much more wide interval of parameters and with much better accuracy. Among them I underline those which look most interesting now.

- Study of photon structure function(s) at small x .
- Nature of growth of **total cross sections**. The widespread concepts assume standard Regge type factorization and universal energy behavior for different processes. With Photon Colliders –*for the first time in particle physics*– one can have the set of mass shell cross sections of very high energy processes, appropriate for the testing of factorization or the level of its violation. They are σ_{pp} , measured at Tevatron and LHC, $\sigma_{\gamma p}$, measured at HERA, $\sigma_{\gamma\gamma}$, measurable at Photon Collider. For this goal, *the preliminary stage of operations with low luminosity can be used* to observe large enough cross sections at small scattering angles.
- In this very low luminosity stage the study of events with particle production in the center of rapidity scale and with rapidity gaps for two vector mesons or dijets originated from initial photons will be crucial for the understanding of **double Pomeron effects** having no reliable explanation till now.
- The study of **charge asymmetry of produced hadrons** in $\gamma\gamma$ collisions will give quite new information about quark–gluon matter at small distances. The charge asymmetry of the produced hadrons in the γe collisions with transverse momentum of scattered electron $p_\perp \geq 30 \text{ GeV}$ will show in explicit form the relation between hadron states produced by vector and axial current.

6 Using of conversion region as γe collider

The conversion region in γe collider with c.m.s. energy about 1.2 MeV but with luminosity about $0.1 \text{ fb}^{-1}/\text{sec}$! It will be unique source of light goldstone particles (axions, majorons, etc.) *LGP*, expected in numerous schemes [10].

The production processes are

$$e\gamma_0 \rightarrow e(\text{LGP}), \quad \gamma\gamma_0 \rightarrow (\text{LGP}). \quad (2)$$

The energy of *LGP* is limited from above as

$$E_{\text{LGP}} \leq \frac{x + a^2 + \sqrt{(x - a^2)^2 - 4a^2}}{2(x + 1)} E, \quad a = \frac{m_a}{m_e}. \quad (3)$$

The angular spread of these *LGP*'s is even narrower than that of high energy photons. In comparison with neutrinos produced by photons in the final wall, the main part of these "axions" is concentrated in a $10^6 \div 10^8$ times more narrow solid angle.

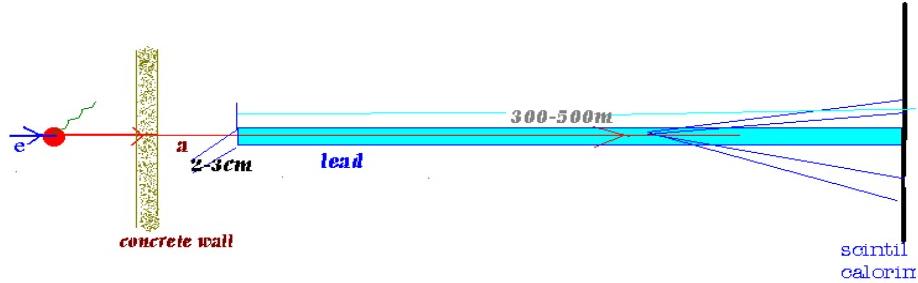


Figure 5: Possible scheme for recording of the "axion"

These *LGP*s interact with matter very weakly. So, the registration scheme can be of the type shown in Fig. 4 (where magnetic bending of residual electrons after conversion is also shown): the *LGP* produced after concrete wall travel to the detector inside $L = 300 - 500$ m long lead wire of diameter 3-5 cm. Here it can interact with hadrons in the nucleus of lead and produce hadrons with typical mean transverse momentum about 300-500 MeV. They can be detected in scintillator-like devices of diameter 3-5 m in the end of wire.

7 Conclusion

This discussion shows that *R&D for Photon collider mode should be performed simultaneously with that of e^+e^- mode of TESLA, etc. Final decision about the turn of different stages of Linear collider should be made only AFTER first operations of LHC*. It can imagine even the opportunity that the Photon collider mode will be switched on before e^+e^- mode. The advantages are:

- The basic electron energy is lower.
- Positron beam is unnecessary.

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